STATISTICAL MOLECULAR THERMODYNAMICS

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Video 4.1

Ideal Monatomic Gas: q_{trans}

IDEAL GAS PARTITION FUNCTION

We have established that the partition function of a *non-interacting* system, Q(N,V,T), can be written in terms of individual atomic or molecular partition functions, q(V,T):

$$Q(N,V,T) = \frac{q(V,T)^N}{N!}$$

N: number of particles

V: volume

T: temperature

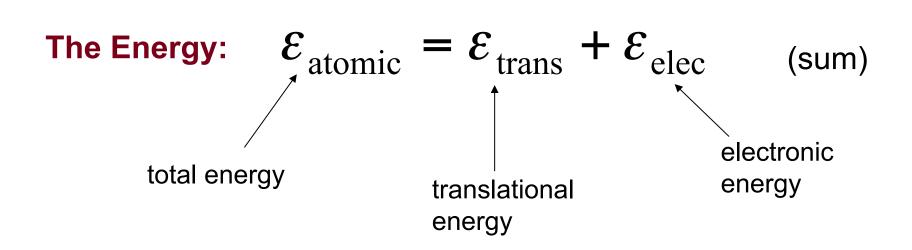
We now apply this to an ideal gas where:

- 1) The molecules are indistinguishable
- 2) The number of states greatly exceeds the number of molecules (assumption of "low" pressure).

MONATOMIC IDEAL GAS PARTITION FUNCTION

We consider each *degree of freedom* separately:

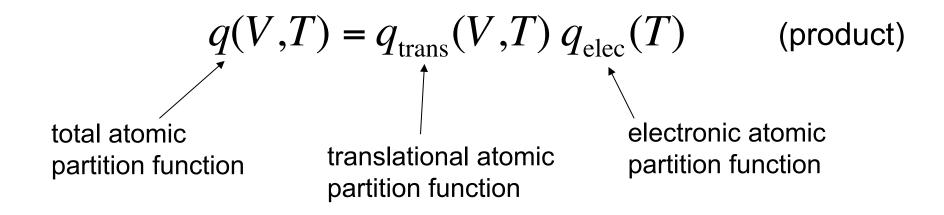
For an ideal *monatomic* gas the *only* degrees of freedom are translational and electronic (a significant simplification):



MONATOMIC IDEAL GAS PARTITION FUNCTION

The atomic partition function is the product of the partition functions from each degree of freedom:

The Partition Function:



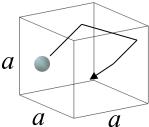
So, we must consider $q_{\rm trans}(V,T)$ and $q_{\rm elec}(T)$

TRANSLATIONAL COMPONENT

Given the general form of the partition function:

$$q_{\rm trans} = \sum_{\rm states} e^{-\beta \varepsilon_{\rm trans}}$$

From QM, we have that for a cubic container with sides of length a,



$$\varepsilon_{\text{trans}} = \frac{h^2}{8ma^2} (n_x^2 + n_y^2 + n_z^2)$$
 with $n_x, n_y, n_z = 1, 2, ...$

For this $\varepsilon_{\rm trans}$

$$q_{\text{trans}} = \sum_{n_x, n_y, n_z = 1}^{\infty} e^{-\beta \varepsilon_{n_x, n_y, n_z}} = \sum_{n_x = 1}^{\infty} \sum_{n_y = 1}^{\infty} \sum_{n_z = 1}^{\infty} \exp \left[-\frac{\beta h^2}{8ma^2} \left(n_x^2 + n_y^2 + n_z^2 \right) \right]$$

TRANSLATIONAL COMPONENT

$$q_{\text{trans}} = \sum_{n_x, n_y, n_z = 1}^{\infty} e^{-\beta \varepsilon_{n_x, n_y, n_z}} = \sum_{n_x = 1}^{\infty} \sum_{n_y = 1}^{\infty} \sum_{n_z = 1}^{\infty} \exp \left[-\frac{\beta h^2}{8ma^2} (n_x^2 + n_y^2 + n_z^2) \right]$$

Or,

$$q_{\text{trans}} = \sum_{n_x=1}^{\infty} \exp\left(-\frac{\beta h^2 n_x^2}{8ma^2}\right) \sum_{n_y=1}^{\infty} \exp\left(-\frac{\beta h^2 n_y^2}{8ma^2}\right) \sum_{n_z=1}^{\infty} \exp\left(-\frac{\beta h^2 n_z^2}{8ma^2}\right)$$

The three sums are all the same (only the index is different), so,

$$q_{\text{trans}}(V,T) = \left[\sum_{n=1}^{\infty} \exp\left(-\frac{\beta h^2 n^2}{8ma^2}\right)\right]^3$$

because of dependence on side length a

DENSE ENERGY LEVELS

There is no simple, analytical expression for this sum:

$$q_{\text{trans}}(V,T) = \left[\sum_{n=1}^{\infty} \exp\left(-\frac{\beta h^2 n^2}{8ma^2}\right)\right]^3$$

However, since *translational energy levels are spaced very densely*, the sum is nearly a continuous function and we can *approximate the sum as an integral*,

$$q_{\text{trans}}(V,T) = \left(\int_{1}^{\infty} dn \, e^{-\beta h^2 n^2 / 8ma^2}\right)^3$$

SOLVING FOR q_{trans}

$$\int_{1}^{\infty} dn \, e^{-\beta h^{2} n^{2} / 8ma^{2}} \approx \int_{0}^{\infty} dn \, e^{-\beta h^{2} n^{2} / 8ma^{2}}$$

Integral table:
$$\int_{0}^{\infty} dn \, e^{-\alpha n^2} = \left(\frac{\pi}{4\alpha}\right)^{1/2}$$

$$q_{\text{trans}}(V,T) \approx \left(\int_{0}^{\infty} dn \, e^{-\beta h^{2} n^{2} / 8ma^{2}}\right)^{3} = \left(\frac{2\pi m k_{\text{B}} T}{h^{2}}\right)^{3/2} V_{\text{N}}$$

Consider methane (CH₄):

$$m = 2.662 \text{ x } 10^{-26} \text{ kg}, V = 24.47 \text{ dm}^3, T = 298.15 \text{ K}$$

$$q_{\text{trans}} = 1.519 \text{ x } 10^{30}$$

So, number of "accessible levels" vastly exceeds Avogradro's number

choice of volume part of "standard-state" convention

 $\alpha = \frac{h^2}{8ma^2k_BT}$