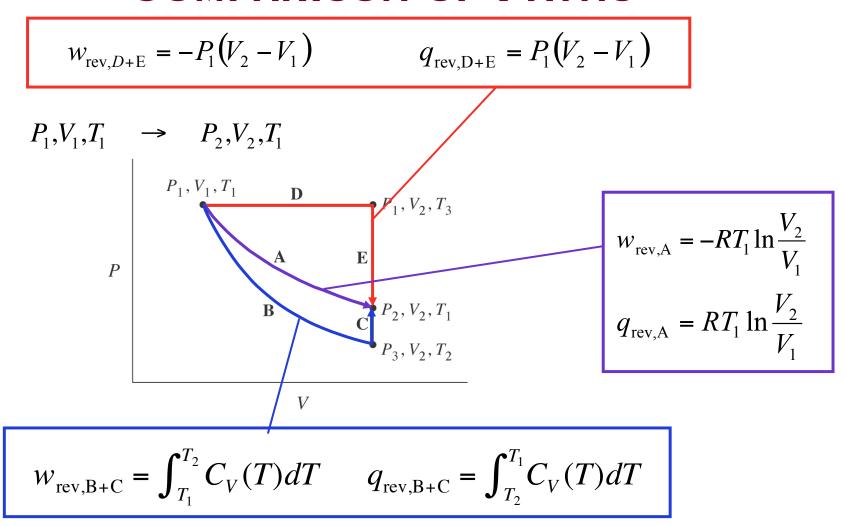
STATISTICAL MOLECULAR THERMODYNAMICS

Christopher J. Cramer

Video 5.5

Adiabatic Processes

COMPARISON OF PATHS



 ΔU =0 for all paths (state function), but $q_{\rm rev}$ and $w_{\rm rev}$ differ

QUANTITATIVE COMPARISON OF PATHS

For example if $P_1 = 4.0$ bar, $V_1 = 0.5$ dm³, $P_2 = 2.0$ bar, $V_2 = 1.0$ dm³, and we have 0.1 moles of ideal monatomic gas:

$$w_{\text{rev},D+E} = -200 \text{ J} \qquad q_{\text{rev},D+E} = 200 \text{ J} \qquad \Delta U_{D+E} = 0$$

$$P_1, V_1, T_1 \rightarrow P_2, V_2, T_1$$

$$P_1, V_1, T_1 \rightarrow P_2, V_2, T_1$$

$$Q_{\text{rev},A} = -139 \text{ J}$$

$$Q_{\text{rev},A} = 139 \text{ J}$$

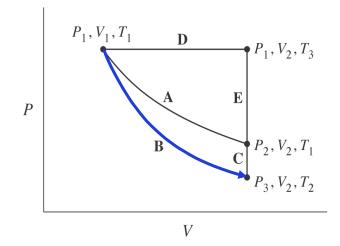
$$\Delta U_A = 0$$

$$v_{\text{rev},B+C} = -111 \text{ J} \qquad q_{\text{rev},B+C} = 111 \text{ J} \qquad \Delta U_{B+C} = 0$$

ADIABATIC EXPANSION COOLS A GAS

Adiabatic, so q = 0 and $dU = \delta w = dw$

(note that if either $\delta q = 0$ or $\delta w = 0$ then the remaining differential becomes exact)



For an ideal gas reversible expansion:

$$dw = dU = C_V(T)dT$$
 and $dw = -PdV = -\frac{nRTdV}{V}$

Putting them together,

$$C_V(T)dT = -\frac{nRT}{V}dV \longrightarrow \int_{T_1}^{T_2} \frac{\overline{C}_V(T)}{T}dT = -R\int_{V_1}^{V_2} \frac{dV}{V} = -R\ln\frac{V_2}{V_1}$$

For a monatomic ideal gas, $\overline{C}_V = \frac{3R}{2}$

$$\frac{3R}{2} \int_{T_1}^{T_2} \frac{dT}{T} = \frac{3R}{2} \ln \frac{T_2}{T_1} = -R \ln \frac{V_2}{V_1} \longrightarrow \left(\frac{T_2}{T_1}\right)^{3/2} = \frac{V_1}{V_2} \quad \text{The gas cools as it expands}$$

ADIABATIC VS ISOTHERMAL IDEAL GAS LAW

Boyle's law for an isothermal process:

$$P_1V_1 = P_2V_2$$

Cf. an adiabatic process (ideal monatomic gas):

$$\left(\frac{T_2}{T_1}\right)^{3/2} = \frac{V_1}{V_2} \xrightarrow{PV = nRT} \left(\frac{P_2 V_2}{P_1 V_1}\right)^{3/2} = \frac{V_1}{V_2}$$

$$P_1 V_1^{5/3} = P_2 V_2^{5/3}$$

less compression; with nowhere to dump heat, temperature rises