

# STATISTICAL MOLECULAR THERMODYNAMICS

*Christopher J. Cramer*

Video 6.6

Entropy and the Partition Function

# ENTROPY: PROBABILITY FORM

Recall from Video 6.4:  $S_{\text{ensemble}} = k_B \left( A \ln A - \sum_j a_j \ln a_j \right)$

where  $A$  is the total number of systems in the ensemble, and  $a_j$  is the population of each system  $j$ .

Then the average system entropy is  $S_{\text{ensemble}}/A$  and the probability  $p_j$  of choosing a system in state  $j$  is  $a_j/A$ ; or,  $a_j = p_j A$

Then:  $S_{\text{ensemble}} = k_B \left( A \ln A - \sum_j p_j A \ln p_j A \right)$

*sum of all probabilities is 1*  
↓

$$= k_B A \ln A - k_B A \sum_j p_j \ln p_j - k_B A \ln A \sum_j p_j$$
$$= -k_B A \sum_j p_j \ln p_j \rightarrow S_{\text{system}} = -k_B \sum_j p_j \ln p_j$$

# ENTROPY: PROBABILITY FORM

$$S_{\text{system}} = -k_B \sum_j p_j \ln p_j$$

Note that L'Hôpital's rule establishes that  $\lim_{x \rightarrow 0} x \ln x = 0$

If all probabilities are 0 except one:  $S = 0$

If all probabilities are equal:  $S$  is maximized

Recall, in the  $NV\beta$  ensemble:  $p_j = \frac{e^{-\beta E_j(N,V)}}{Q(N,V,\beta)}$

$$\rightarrow S = -k_B \sum_j \frac{e^{-\beta E_j}}{Q} \ln \left( \frac{e^{-\beta E_j}}{Q} \right) = -k_B \sum_j \frac{e^{-\beta E_j}}{Q} (-\beta E_j - \ln Q)$$

# ENTROPY AND THE PARTITION FUNCTION

$$S = -k_B \sum_j \frac{e^{-\beta E_j}}{Q} (-\beta E_j - \ln Q)$$

Some manipulation:

$$S = \frac{1}{T} \sum_j p_j E_j + \frac{k_B \ln Q}{Q} \sum_j e^{-\beta E_j}$$

sum is  $Q$

$$= \frac{U}{T} + k_B \ln Q$$

probability weighted energy  
is internal energy  $U$

$$= k_B T \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} + k_B \ln Q$$

$S$  can be computed directly from partition function!

# ENTROPY OF A MONATOMIC IDEAL GAS

$$Q = \frac{1}{N!} \left( \frac{2\pi m k_B T}{h^2} \right)^{3N/2} V^N g_{e1}^N \quad \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} = \frac{3N}{2} \left( \frac{1}{T} \right)$$

$$\ln Q = N \ln \left[ \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} V g_{e1} \right] - \ln N!$$

$$= N \ln \left[ \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} V g_{e1} \right] - N \ln N + N$$

$$= N \ln \left[ \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} \frac{V}{N} g_{e1} \right] + N$$

$$S = k_B T \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} + k_B \ln Q$$

## ENTROPY OF A MONATOMIC IDEAL GAS

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$$\ln Q = N \ln \left[ \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} \frac{V}{N} g_{e1} \right] + N \quad S = k_B T \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} + k_B \ln Q$$

If  $N = N_A$  (molar quantity):

$$\bar{S} = \frac{5}{2} R + R \ln \left[ \left( \frac{2\pi m k_B T}{h^2} \right)^{3/2} \frac{V g_{e1}}{N_A} \right]$$

Entropy increases with 1) increasing mass  $m$ , 2) increasing temperature  $T$ , 3) increasing standard-state volume  $V$ , and 4) increasing electronic ground-state degeneracy  $g_{e1}$