

STATISTICAL MOLECULAR THERMODYNAMICS

Christopher J. Cramer

Video 9.7

Chemical Potential from the Partition Function

CHEMICAL POTENTIAL DEFINED FOR A

Just as with the Gibbs free energy, we can take the total differential of the Helmholtz free energy with moles of substance taken as a natural variable, i.e.,

$$\begin{aligned}dA &= \left(\frac{\partial A}{\partial T}\right)_{n,V} dT + \left(\frac{\partial A}{\partial V}\right)_{n,T} dV + \left(\frac{\partial A}{\partial n}\right)_{T,V} dn \\ &= -SdT - PdV + \left(\frac{\partial A}{\partial n}\right)_{T,V} dn\end{aligned}$$

Now, recalling that $G = A + PV$, we determine dG as

$$dG = dA + d(PV) = -SdT - \cancel{PdV} + \left(\frac{\partial A}{\partial n}\right)_{T,V} dn + \cancel{PdV} + VdP$$

EQUALITY OF CHEMICAL POTENTIALS

$$dG = -SdT + VdP + \left(\frac{\partial A}{\partial n}\right)_{T,V} dn$$

But, direct expansion of the total differential of G gives

$$dG = \left(\frac{\partial G}{\partial T}\right)_{P,n} dT + \left(\frac{\partial G}{\partial P}\right)_{T,n} dP + \left(\frac{\partial G}{\partial n}\right)_{T,P} dn$$

$$dG = -SdT + VdP + \mu dn$$

so evidently, the two chemical potentials are equal when evaluated holding the respective natural variables constant

$$\mu = \left(\frac{\partial G}{\partial n}\right)_{T,P} = \left(\frac{\partial A}{\partial n}\right)_{T,V}$$

WHAT MAKES A SO USEFUL?

Recall the equations for U and S as functions of Q

$$\begin{aligned} U &= k_B T^2 \left(\frac{\partial \ln Q}{\partial T} \right)_{N,V} \\ S &= k_B T \left(\frac{\partial \ln Q}{\partial T} \right)_{N,V} + k_B \ln Q \end{aligned} \quad \longrightarrow \quad \begin{aligned} A &= U - TS \\ A &= -k_B T \ln Q \end{aligned}$$

This offers a very direct way to compute the chemical potential from the partition function, based on the latter's dependence on number of particles (or moles)!

RELATING CHEMICAL POTENTIAL TO Q

$$\mu = \left(\frac{\partial G}{\partial n} \right)_{T,P} = \left(\frac{\partial A}{\partial n} \right)_{T,V} = -k_B T \left(\frac{\partial \ln Q}{\partial n} \right)_{V,T} = -RT \left(\frac{\partial \ln Q}{\partial N} \right)_{V,T}$$

remember:

$$Q(N, V, T) = \frac{q(V, T)^N}{N!} \quad \text{so} \quad \ln Q = N \ln q - N \ln N + N$$

*Stirling's approximation
used for $\ln(N!)$*


Thus,

$$\mu = -RT (\ln q - \ln N - 1 + 1) = -RT \ln \frac{q(V, T)}{N}$$

RELATING CHEMICAL POTENTIAL TO Q

$$\mu = -RT \ln \frac{q(V, T)}{N}$$

Recall q is linear in V and write:

$$\begin{aligned} \mu &= -RT \ln \left[\frac{q(T)}{V} \cdot \frac{V}{N} \right] \\ &= -RT \ln \left[\frac{q(T)}{V} \cdot \frac{k_B T}{P} \right] \\ &= -RT \ln \left[\frac{q(T)}{V} \cdot k_B T \right] + RT \ln P \end{aligned}$$


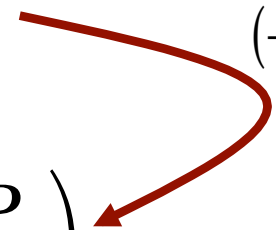
ideal gas

USING STANDARD STATE CONVENTION

Pressure is expressed relative to standard state (1 bar). We emphasize by writing:

$$\mu(T, P) = -RT \ln \left[\left(\frac{q}{V} \right) k_B T \right] + RT \ln P$$

Add zero
($-RT \ln P^\circ + RT \ln P^\circ$)



$$\mu(T, P) = -RT \ln \left[\left(\frac{q}{V} \right) \frac{k_B T}{P^\circ} \right] + RT \ln \left(\frac{P}{P^\circ} \right)$$

$$\mu(T, P) = \mu^\circ(T) + RT \ln \left(\frac{P}{P^\circ} \right)$$

cf. from slide 1 of video 8.7: $\bar{G}(T, P) = G^\circ(T) + RT \ln P$

chemical potential more general — and connects other phases to gas phase!!

$$dU = \delta q + \delta w$$



Next: Review of Module 9